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## Limit cycles of a second-order differential equation

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ABSTRACT

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## 1. Introduction

To determine the number of limit cycles of a differential equation is one of the main problems in the qualitative theory of planar differential system. In 1881 Poincaré [1] defined the notion of *limit cycle* of a planar differential system as a periodic orbit isolated in the set of all periodic orbits of the differential system. And he defined the notion of a *center* of a real planar differential system, i.e. of an isolated equilibrium point having a neighborhood filled with periodic orbits. Later on one way to produce limit cycles is by perturbing the periodic orbits of a center, see for instance the papers [2–5] and the references quoted there.

In [6] Mathieu considered the second order differential equation

$$\ddot{x} + b(1 + \cos t)x = 0, \tag{1}$$

where b is a real constant. It is called Mathieu equation, which is the simplest mathematical model of an excited system depending on a parameter. The more general Ermakov–Pinney equation is the Mathieu–Duffing type equation

$$\ddot{x} + b(1 + \cos t)x - x^{\beta} = 0, \tag{2}$$

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where  $\varepsilon > 0$  is a small parameter, *m* is an arbitrary non-negative integer, Q(x, y) is a polynomial of degree *n* and  $\theta = \arctan(y/x)$ . The main tool used for proving our

We provide an upper bound for the maximum number of limit cycles bifurcating

from the periodic solutions of  $\ddot{x} + x = 0$ , when we perturb this system as follows

 $\ddot{x} + \varepsilon (1 + \cos^m \theta) Q(x, y) + x = 0,$ 

results is the averaging theory.

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