On the dynamics of an affine system with non-reducible linear behaviour



Marc Jorba, Àngel Jorba Universitat de Barcelona

Description of the problem

Consider the following quasi-periodically forced discrete dynamical system:

 $\begin{cases} \bar{\mathbf{x}} = f(\mathbf{x}, \theta), \\ \bar{\theta} = \theta + \omega. \end{cases}$

Where $\mathbf{x} \in \mathbf{U}, \theta \in \mathbb{T}, \mathbf{f} : \mathbf{U} \times \mathbb{T} \mapsto \mathbf{U}$ is of class $\mathbf{C}^{r+1}, \mathbf{r} \ge \mathbf{0}$, and $\omega \in (0, 2\pi) \setminus 2\pi \mathbb{Q}$. Suppose there exists an invariant curve, a function $\boldsymbol{x} \in \boldsymbol{x}$

 $C^{r}(\mathbb{T}, U)$, such that

 $\mathbf{x}(\theta + \omega) = \mathbf{f}(\mathbf{x}(\theta), \theta).$

The linear dynamics around the **invariant curve** is described by:

 $\begin{cases} \bar{\mathbf{x}} = \mathbf{A}(\theta)\mathbf{x}, \\ \bar{\theta} = \theta + \omega. \end{cases}$

We have named $A(\theta) := D_X f(x(\theta))$. Any system such as (1) is called quasi-periodic linear skew product, or a quasi-periodic *C^r* cocycle.

Reducibility

Essential non-reducibility

Reducibility can have a topological obstruction. In that case we speak about essential non-reducibility:

Definition

(1)

Let $A \in C^r(\mathbb{T}, \operatorname{GL}_2 \mathbb{R})$, $r \geq 0$. Fix a vector $v \in \mathbb{R}^2 \setminus \{0\}$ and consider the curve v_A at $\mathbb{R}^2 \setminus \{0\}$ given by $v_A(\theta) = A(\theta)v$. We define the winding number of **A**, wind **A**, as the winding number of v_A around the origin of \mathbb{R}^2 .

- The winding number of a cocycle does not depend on the choice of the vector **v**.
- ▶ The winding number is invariant under conjugation. Since a constant matrix **B** verifies wind B = 0, if **A** is a cocycle such that wind $A \neq 0$, then A is essentially non-reducible.
- A linear map coming from a Poincaré section of a quasiperiodic linear ODE has winding number zero.

Goal of this work

► We show that non-reducibility has **dynamical manifestation**.

We study the destruction of the family of attracting invariant curves of the model:

$$\begin{cases} \left(\frac{\bar{\boldsymbol{x}}}{\bar{\boldsymbol{y}}} \right) &= \mu \left(\begin{array}{c} \cos \theta &- \sin \theta \\ \sin \theta & \cos \theta \end{array} \right) \left(\begin{array}{c} \boldsymbol{x} \\ \boldsymbol{y} \end{array} \right) + \left(\begin{array}{c} \boldsymbol{v_1} \\ \boldsymbol{v_2} \end{array} \right), \quad (2) \\ \bar{\theta} &= \theta + \omega. \end{cases}$$

• Here, μ is a **positive** parameter.

- ► The linear behaviour of the model is essentially non-reducible.
- ▶ There exists a smooth invariant curve z_{μ} for each $\mu \neq 1$.
- The curve z_{μ} is attracting when $\mu < 1$.

Theorem

Consider the system (2). Assume the rotation number ω to be of constant type. Then, when $\mu \rightarrow 1$:

1. The invariant curve \mathbf{z}_{μ} undergoes a **fractalization** process, i.e.

$$\frac{\|\boldsymbol{z}_{\boldsymbol{\mu}}'\|_{\infty}}{\|\boldsymbol{z}_{\boldsymbol{\mu}}\|_{\infty}} = \mathcal{O}(1-\boldsymbol{\mu})^{-1}.$$

Definition

System (1) is said to be C^{r} -reducible if there exists a C^{r} change of variables $\mathbf{x} = \mathbf{C}(\theta)\mathbf{y}$ such that transforms the former system into:

 $\begin{cases} \bar{\mathbf{y}} = \mathbf{B}\mathbf{y}, \\ \bar{\theta} = \theta + \omega. \end{cases}$

Where the matrix $\mathbf{B} = \mathbf{C}^{-1}(\theta + \omega)\mathbf{A}(\theta)\mathbf{C}(\theta)$ does not depend on θ .

Affine systems Consider the following affine system on the plane:

 $\begin{cases} \bar{\mathbf{x}} = \mu \mathbf{A}(\theta) \mathbf{x} + \mathbf{b}(\theta), \\ \bar{\theta} = \theta + \omega, \end{cases}$

with $A \in C^r(\mathbb{T}, \operatorname{GL}_2\mathbb{R})$ and $b \in C^r(\mathbb{T}, \mathbb{R}^2) =: E, r \ge 0$, endowed with the standard C^r norm. Let $||A|| = \sup_{||x||=1} ||Ax||$

and $\rho(\mathbf{A}) = \lim_{k \to \infty} \|\mathbf{A}^k\|^{1/k}$.

• If $\rho(A) < 1$ an attracting invariant curve appears as a fixed point of the operator:

 $\mathcal{T}(\mathbf{x}(\theta)) = \mu \mathbf{A}(\theta - \omega)\mathbf{x}(\theta - \omega) + \mathbf{b}(\theta - \omega).$

 $\|-\mu\|\infty$

2. The winding number of \mathbf{z}_{μ} around any point of \mathbb{R}^2 verifies the following:

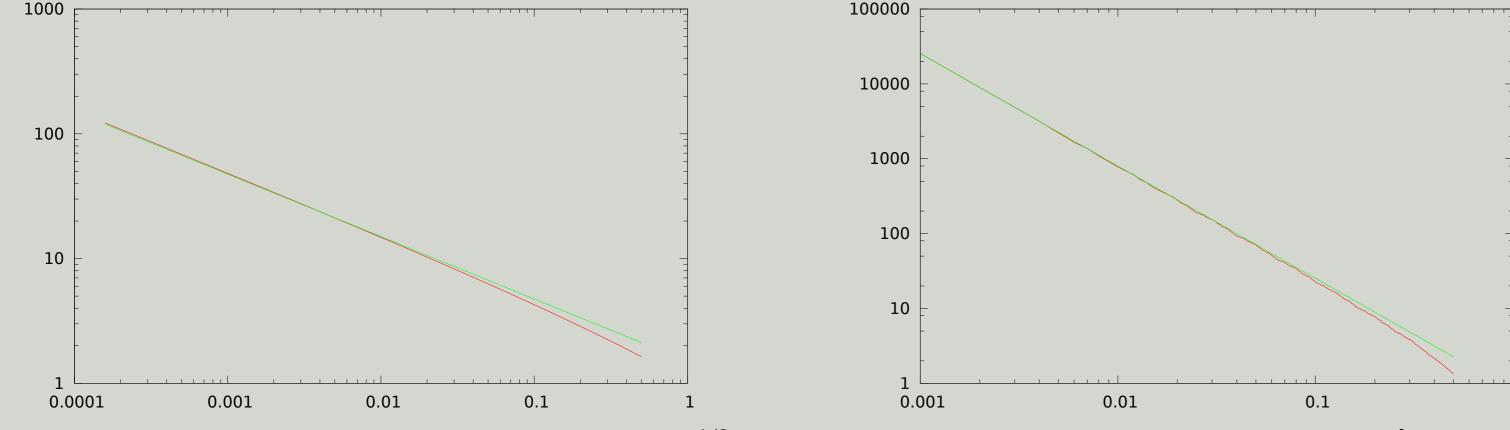
wind $z_{\mu} = O(1 - \mu)^{-1}$.

Corollary

Let \mathbf{z}_{μ} be the invariant curve of (2). It holds

$$\bigcup_{\mu \in (\mu_0, 1)} \operatorname{graph} z_{\mu} = \mathbb{R}^2 \quad \text{for any } \mu_0 \in (0, 1).$$

Numerical experiments: exploring the breakdown scenario



Sketch the proof

We take advantage on the identification between \mathbb{R}^2 and the complex plane.

 $\begin{cases} \bar{z} = \mu e^{i\theta} z + c, \\ \bar{\theta} = \theta + \omega. \end{cases}$

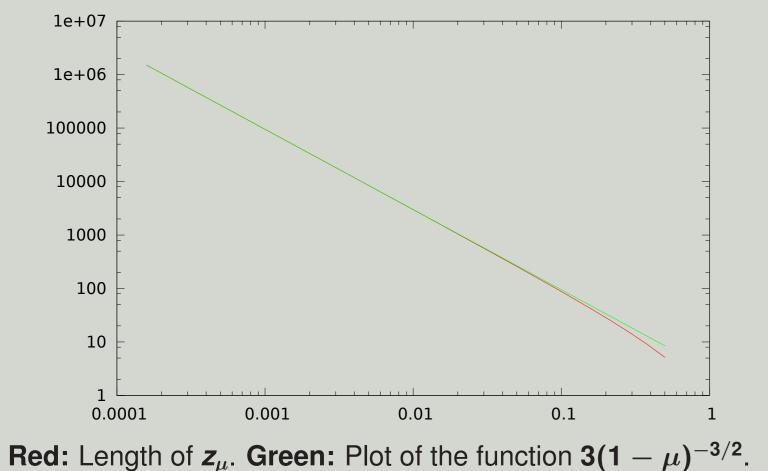
It is possible to compute explicitly the Fourier expansion for the invariant curve:

$$\boldsymbol{z}_{\mu}(\boldsymbol{\theta}) = \boldsymbol{c} \sum_{\boldsymbol{k}=0}^{\infty} \mu^{\boldsymbol{k}} \boldsymbol{e}^{-\boldsymbol{i}\frac{\boldsymbol{k}(\boldsymbol{k}+1)}{2}\omega} \boldsymbol{e}^{\boldsymbol{i}\boldsymbol{k}\boldsymbol{\theta}}, \qquad (3)$$

which is convergent whenever $\mu < 1$. The series was studied by Hardy and Littlewood in 1914, from their work it is easy to see that, if ω is of **constant type**:

$$\| oldsymbol{z}_{\mu} \|_{\infty} = \mathcal{O}(1-\mu)^{-1/2}. \ \| oldsymbol{z}'_{\mu} \|_{\infty} = \mathcal{O}(1-\mu)^{-3/2}.$$

Red: Plot of $||z_{\mu}||_{\infty}$. Green: Plot of the function $2(1-\mu)^{-1/2}$. Red: Plot of $||z'_{\mu}||_{\infty}$. Green: Plot of the function $\frac{4}{5}(1-\mu)^{-3/2}$.

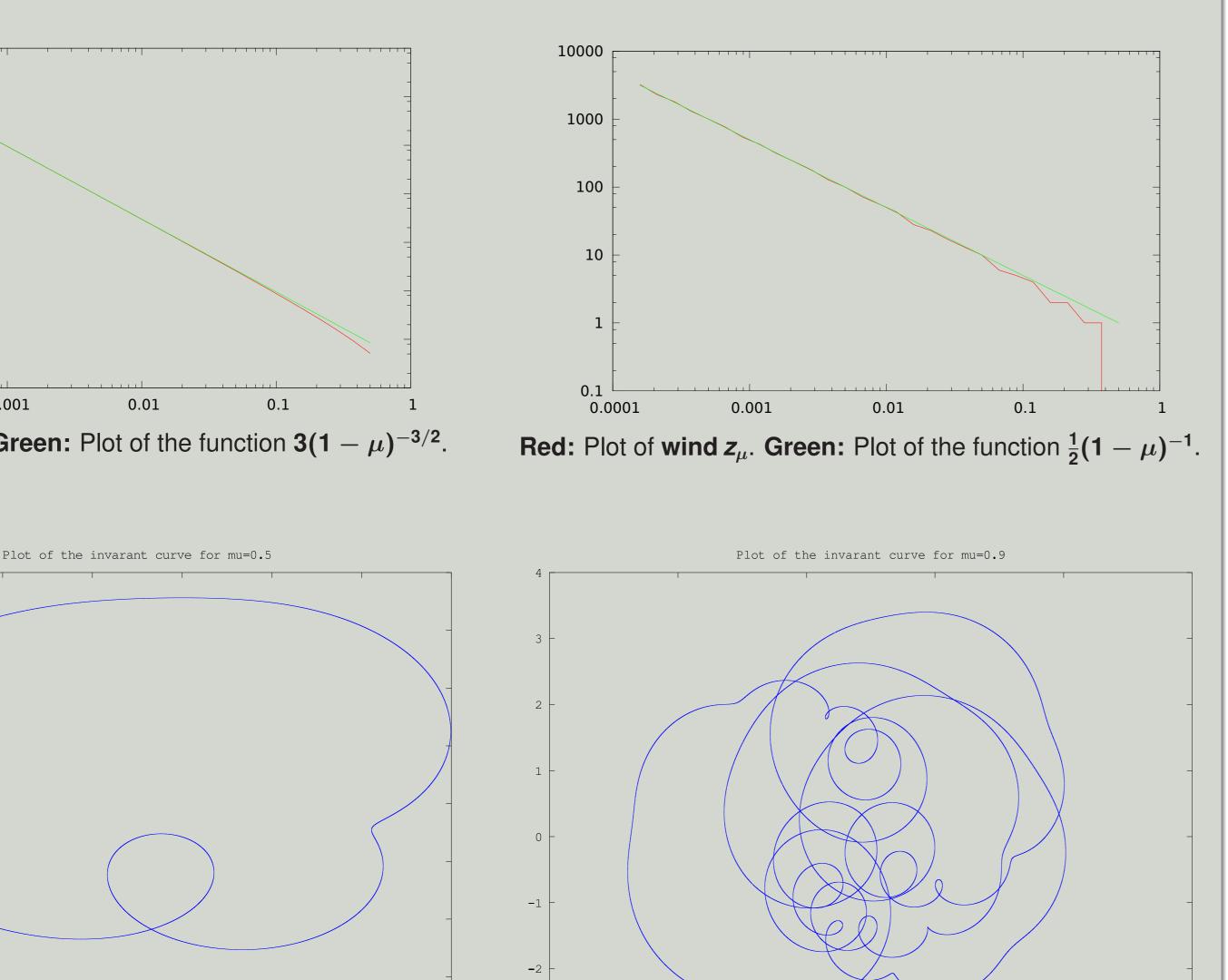


-0.2

-0.4

-0.6

-0.8



By the principle argument, the winding number of (3) is the number of zeros in the unit disk of the function

$$f(z) = c \sum_{k=0}^{\infty} e^{-i \frac{k(k+1)}{2} \omega} z^k.$$

The winding number, hence, can be defined for a full measure set of μ 's and

wind $z_{\mu} = O(1 - \mu)^{-1}$.

On the reducible case

Assume we deal with a reduced affine system with one dimensional complex coordinates:

$$ar{\zeta} = \mu \zeta + \boldsymbol{b}(\theta), \ ar{ heta} = heta + \omega.$$

Assume:

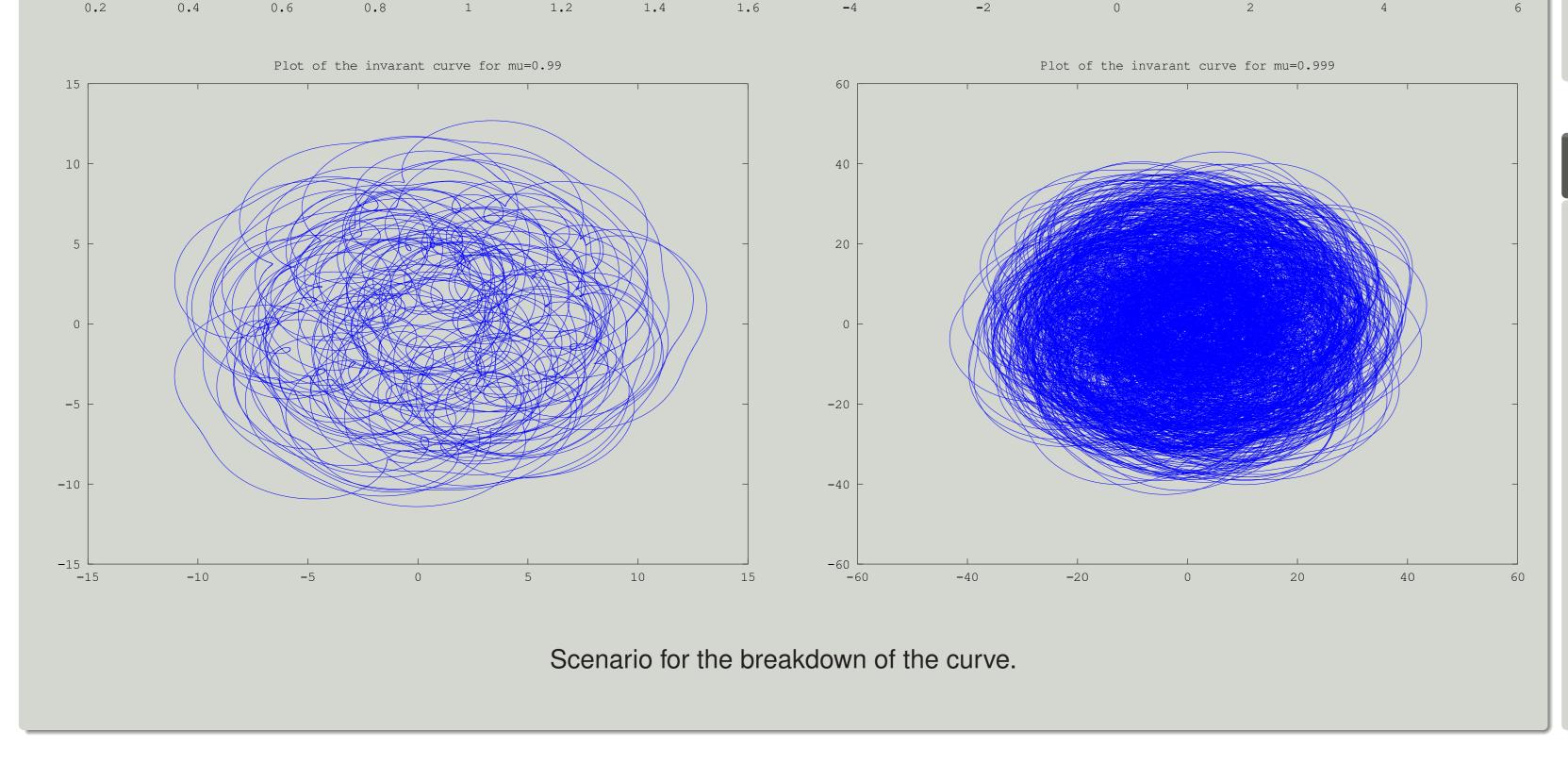
• The rotation number ω is **Diophantine**.

► The independent term **b** is smooth.

We look for a curve $\zeta_{\mu} \in C^{r}(\mathbb{T}, \mathbb{C})$ which satisfies the condition of invariant curve. It is possible to find explicitly the Fourier coefficients.

$$\zeta_{\mathbf{k}}^{\mu} = rac{\mathbf{b}_{\mathbf{k}}}{\mathbf{e}^{i\mathbf{k}\omega} - \mu}, \quad \mathbf{k} \in \mathbb{Z}.$$

The smoothness of **b** implies a suitable decay on the values $|\mathbf{b}_{\mathbf{k}}|$, then: • If **b** is average free, the decay on the values $|\mathbf{b}_{\mathbf{k}}|$, imply the smoothness of ζ_{μ} . • If **b** is average different from zero, the invariant curve diverges when $\mu \to \mathbf{1}$ with velocity $\mathcal{O}(1-\mu)^{-1}$ and with bounded derivative.



No fractalization occurs in the reducible case.

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