

On periodic solutions of 2-periodic Lyness difference equations

Guy Bastien¹, Víctor Mañosa² and Marc Rogalski³

¹Institut Mathématique de Jussieu, Université Paris 6 and CNRS,

²DMA3-CoDALab, Universitat Politècnica de Catalunya*.

³Laboratoire Paul Painlevé, Université de Lille 1; Université Paris 6 and CNRS,

18th International Conference on Difference Equations and Applications
July 2012, Barcelona, Spain.

*Supported by MCYT's grant DPI2011-25822 and SGR program.

We study the *set of periods* of the 2-periodic Lyness' equations

$$u_{n+2} = \frac{a_n + u_{n+1}}{u_n}, \quad (1)$$

where

$$a_n = \begin{cases} a & \text{for } n = 2\ell + 1, \\ b & \text{for } n = 2\ell, \end{cases} \quad (2)$$

and being $(u_1, u_2) \in \mathcal{Q}^+$; $\ell \in \mathbb{N}$ and $a > 0, b > 0$.

This can be done using the *composition map*:

$$F_{b,a}(x, y) := (F_b \circ F_a)(x, y) = \left(\frac{a+y}{x}, \frac{a+bx+y}{xy} \right), \quad (3)$$

where F_a and F_b are the Lyness maps: $F_\alpha(x, y) = (y, \frac{\alpha+y}{x})$. Indeed:

$$(u_1, u_2) \xrightarrow{F_a} (u_2, u_3) \xrightarrow{F_b} (u_3, u_4) \xrightarrow{F_a} (u_4, u_5) \xrightarrow{F_b} (u_5, u_6) \xrightarrow{F_a} \dots$$

The map $F_{b,a}$:

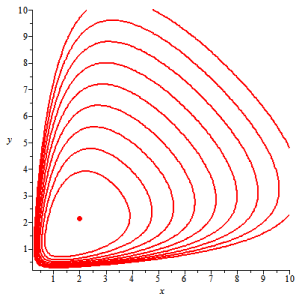
- Is a **QRT** map whose first integral is (Quispel, Roberts, Thompson; 1989):

$$V_{b,a}(x, y) = \frac{(bx + a)(ay + b)(ax + by + ab)}{xy},$$

see also (Janowski, Kulenović, Nurkanović; 2007) and (Feuer, Janowski, Ladas; 1996).

- Has a **unique** fixed point $(x_c, y_c) \in \mathcal{Q}^+$, which is the **unique global minimum** of $V_{b,a}$ in \mathcal{Q}^+ .
- Setting $h_c := V_{b,a}(x_c, y_c)$, for $h > h_c$ the level sets $\{V_{b,a} = h\} \cap \mathcal{Q}^+$ are the **closed curves**.

$$\mathcal{C}_h^+ := \{(bx + a)(ay + b)(ax + by + ab) - hxy = 0\} \cap \mathcal{Q}^+ \text{ for } h > h_c.$$



The dynamics of $F_{b,a}$ restricted to \mathcal{C}_h^+ is *conjugate to a rotation* with associated *rotation number* $\theta_{b,a}(h)$.

Theorem A

Consider the family $F_{b,a}$ with $a, b > 0$.

- (i) If $(a, b) \neq (1, 1)$, then $\exists p_0(a, b) \in \mathbb{N}$ s.t. for any $p > p_0(a, b)$, \exists at least an oval C_h^+ filled by p -periodic orbits.
- (ii) The set of periods arising in the family $\{F_{b,a}, a > 0, b > 0\}$ restricted to \mathbb{Q}^+ contains all prime periods except 2, 3, 4, 6, 10.

Corollary.

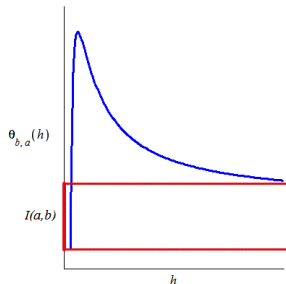
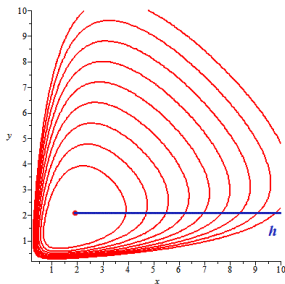
Consider the 2-periodic Lyness' recurrence for $a, b > 0$ and positive initial conditions u_1 and u_2 .

- (i) If $(a, b) \neq (1, 1)$, then $\exists p_0(a, b) \in \mathbb{N}$, s.t. for any $p > p_0(a, b)$ \exists continua of initial conditions giving $2p$ -periodic sequences.
- (ii) The set of prime periods arising when $(a, b) \in (0, \infty)^2$ and positive initial conditions are considered contains all the even numbers except 4, 6, 8, 12, 20.
If $a \neq b$, then it does not appear any odd period, except 1.

The value $p_0(a, b)$ is computable for an open and dense set in the parameter space.

To compute the allowed periods, the main issues to take into account are:

- The fact that the rotation number function $\theta_{b,a}(h)$ is **continuous** in $[h_c, +\infty)$.
- The fact that *generically* $\theta_{b,a}(h_c) \neq \lim_{h \rightarrow +\infty} \theta_{b,a}(h) \Rightarrow \exists I(a, b)$, a **rotation interval**.



Proposition B.

$$\lim_{h \rightarrow h_c^+} \theta_{b,a}(h) = \sigma(a, b) := \frac{1}{2\pi} \arccos \left(\frac{1}{2} \left[-2 + \frac{1}{x_c y_c} \right] \right), \quad \text{and} \quad \lim_{h \rightarrow +\infty} \theta_{b,a}(h) = \frac{2}{5}.$$

Corollary

$$\text{Set } I(a, b) := \left\langle \sigma(a, b), \frac{2}{5} \right\rangle.$$

- If $\sigma(a, b) \neq 2/5 \forall \theta \in I(a, b)$, \exists an oval C_h^+ s.t. $F_{b,a}(C_h^+)$ is *conjugate to a rotation*, with a rotation number $\theta_{b,a}(h) = \theta$.
- In particular, \forall irreducible $q/p \in I(a, b)$, \exists *periodic orbits* of $F_{b,a}$ of **prime period** p .

The periods of the family $F_{b,a}$.

Using the previous results with the family $a = b^2$ we found that:

$$\bigcup_{b>0} I(b^2, b) = \left(\frac{1}{3}, \frac{1}{2}\right) \subset \bigcup_{a>0, b>0} I(a, b) \subset \bigcup_{a>0, b>0} \text{Image}(\theta_{b,a}(h_c, +\infty)).$$

Proposition.

- For each θ in $(1/3, 1/2) \exists a, b > 0$ and an oval C_h^+ , s.t. $F_{b,a}(C_h^+)$ is conjugate to a rotation with rotation number $\theta_{b,a}(h) = \theta$.
- In particular, \forall irreducible $q/p \in (1/3, 1/2), \exists p$ -periodic orbits of $F_{b,a}$

We'll know some periods of $\{F_{b,a}, a, b > 0\}$

\Leftrightarrow

We know which are the irreducible fractions in $(1/3, 1/2)$

Lemma (Cima, Gasull, M; 2007)

Given (c, d) ; Let $p_1 = 2, p_2 = 3, p_3, \dots, p_n, \dots$ be all the prime numbers.

- Let p_{m+1} be the smallest prime number satisfying that $p_{m+1} > \max(3/(d - c), 2)$,
- Given any prime number $p_n, 1 \leq n \leq m$, let s_n be the smallest natural number such that $p_n^{s_n} > 4/(d - c)$.
- Set $p_0 := p_1^{s_1-1} p_2^{s_2-1} \dots p_m^{s_m-1}$.

Then, for any $p > p_0 \exists$ an irreducible fraction q/p s.t. $q/p \in (c, d)$.

Proof of Theorem A (ii):

- We apply the above result to $(1/3, 1/2)$. $\forall p \in \mathbb{N}$, s.t. $p > p_0$

$$p_0 := 2^4 \cdot 3^3 \cdot 5 \cdot 7 \cdot 11 \cdot 13 \cdot 17 = 12\,252\,240,$$

\exists an irreducible fraction $q/p \in (1/3, 1/2)$.

- A finite checking determines which values of $p \leq p_0$ s.t. $q/p \in (1/3, 1/2)$, resulting that there appear irreducible fractions with all the denominators except **2, 3, 4, 6 and 10**.
- Proposition C $\implies \exists a, b > 0$ s.t. \exists an oval with rotation number $\theta_{b,a}(h) = q/p$, thus giving rise to p -periodic orbits of $F_{b,a}$ for all allowed p .
- Still it must be proved that **2, 3, 4, 6 and 10** are forbidden, since

$$I(a, b) \subseteq \text{Image}(\theta_{b,a}(h_c, +\infty))$$

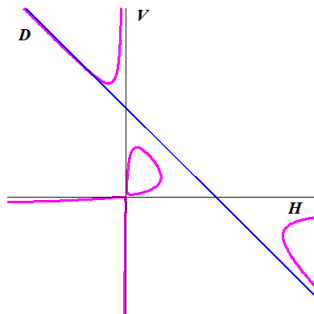


Continuity and asymptotic behavior of $\theta_{b,a}(h)$.

The curves \mathcal{C}_h , in homogeneous coordinates $[x : y : t] \in \mathbb{C}P^2$, are

$$\tilde{\mathcal{C}}_h = \{(bx + at)(ay + bt)(ax + by + abt) - hxyt = 0\}.$$

The points $H = [1 : 0 : 0]$; $V = [0 : 1 : 0]$; $D = [b : -a : 0]$ are common to all curves



Proposition

If $a > 0$ and $b > 0$, and for all $h > h_c$, the curves $\tilde{\mathcal{C}}_h$ are *elliptic*.

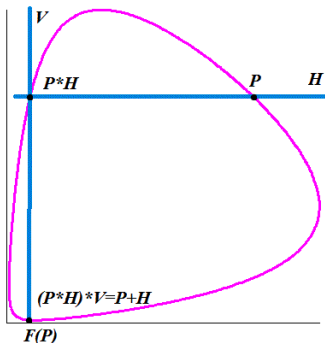
$F_{b,a}$ extends to $\mathbb{C}P^2$ as $\tilde{F}_{b,a}([x : y : t]) = [ayt + y^2 : at^2 + bxt + yt : xy]$.

Lemma. Relation between the dynamics of $F_{b,a}$ and the group structure of \mathcal{C}_h (*)

For each h s.t. $\tilde{\mathcal{C}}_h$ is elliptic,

$$\tilde{F}_{b,a}|_{\tilde{\mathcal{C}}_h}(P) = P + H$$

Where $+$ is the addition of the group law of $\tilde{\mathcal{C}}_h$ taking the infinite point V as the zero element.



Observe that

$$F^n(P) = P + nH,$$

so $\tilde{\mathcal{C}}_h$ is full of *p-periodic orbits* \Leftrightarrow

$$pH = V$$

i.e. H is a *torsion point* of $\tilde{\mathcal{C}}_h$.

(*) Birational maps preserving elliptic curves can be explained using its group structure (Jogia, Roberts, Vivaldi; 2006).

Instead of looking to a normal form for F we look for a normal form for $\tilde{\mathcal{C}}_h$.

$$\begin{aligned} \left(\tilde{\mathcal{C}}_h, +, V \right) & \xrightarrow{\cong} \left(\hat{\mathcal{E}}_L, +, \hat{V} \right) \\ \tilde{F}|_{\tilde{\mathcal{C}}_h} : P \mapsto P + H & \longrightarrow \hat{G}|_{\mathcal{E}_L} : P \mapsto P + \hat{H} \end{aligned}$$

Where $\hat{\mathcal{E}}_L$ is the *Weierstrass Normal Form* which in the affine plane is:

$$\mathcal{E}_L = \{ y^2 = 4x^3 - g_2x - g_3 \}$$

with $g_i := g_i(a, b, h)$.

WHY?

- 1 Because we can *parameterize* it using the Weierstrass \wp function...
- 2 ...that gives an *integral expression* for the *rotation number function*.

$$2\Theta(L) = \frac{\int_{x(L)}^{+\infty} \frac{ds}{\sqrt{4s^3 - g_2s - g_3}}}{\int_{e_1}^{+\infty} \frac{ds}{\sqrt{4s^3 - g_2s - g_3}}} \quad \text{where} \quad \theta_{b,a}(h) \sim \Theta(L)$$

- 3 The *asymptotics* of this integral expression can be studied.

This scheme was used in (Bastien, Rogalski; 2004).

The Weierstrass normal form of C_h is

$$\mathcal{E}_L = \{ y^2 = 4x^3 - g_2x - g_3 \}$$

where

$$g_2 = \frac{1}{192} \left(L^8 + \sum_{i=4}^7 p_i(\alpha, \beta) L^i \right) \quad \text{and} \quad g_3 = \frac{1}{13824} \left(-L^{12} + \sum_{i=6}^{11} q_i(\alpha, \beta) L^i \right),$$

being

$$\begin{aligned} p_7(a, b) &= -4(\alpha + \beta + 1), \\ p_6(a, b) &= 2(3(\alpha - \beta)^2 + 2(\alpha + \beta) + 3), \\ p_5(a, b) &= -4(\alpha + \beta - 1)(\alpha^2 - 4\beta\alpha + \beta^2 - 1), \\ p_4(a, b) &= (\alpha + \beta - 1)^4. \end{aligned}$$

and

$$\begin{aligned} q_{11}(a, b) &= 6(\alpha + \beta + 1), \\ q_{10}(a, b) &= 3(-5\alpha^2 + 2\alpha\beta - 5\beta^2 - 6\alpha - 6\beta - 5) \\ q_9(a, b) &= 4(5\alpha^3 - 12\alpha^2\beta - 12\alpha\beta^2 + 5\beta^3 + 3\alpha^2 - 3\alpha\beta + 3\beta^2 + 3\alpha + 3\beta + 5) \\ q_8(a, b) &= 3(-5\alpha^4 + 16\alpha^3\beta - 30\alpha^2\beta^2 + 16\alpha\beta^3 - 5\beta^4 + 4\alpha^3 \\ &\quad - 12\alpha^2\beta - 12\alpha\beta^2 + 4\beta^3 + 2\alpha^2 - 8\alpha\beta + 2\beta^2 + 4\alpha + 4\beta - 5) \\ q_7(a, b) &= 6(\alpha^2 - 4\alpha\beta + \beta^2 - 1)(\alpha + \beta - 1)^3 \\ q_6(a, b) &= -(\alpha + \beta - 1)^6 \end{aligned}$$

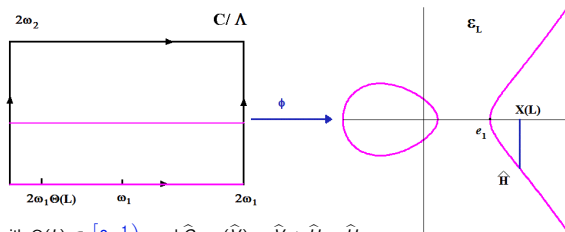
where $\alpha = a/b^2$ and b/a^2 and $L \rightarrow +\infty \Leftrightarrow h \rightarrow +\infty$.

Since $\widehat{\mathcal{E}}_L \cong \mathbb{T}^2 = \mathbb{C}/\Lambda$, the Weierstrass \wp function relative to a lattice Λ gives the parametrization of $\widehat{\mathcal{E}}_L$.

$$\phi : \begin{array}{l} \mathbb{T}^2 = \mathbb{C}/\Lambda \longrightarrow \widehat{\mathcal{E}}_L \\ z \longrightarrow \left\{ \begin{array}{l} [\wp(z) : \wp'(z) : 1] \text{ if } z \notin \Lambda; \\ [0 : 1 : 0] = \widehat{V} \text{ if } z \in \Lambda, \end{array} \right. \end{array}$$

Hence $\wp'(z)^2 = 4\wp(z)^3 - g_2\wp(z) - g_3$ since $y^2 = 4x^3 - g_2x - g_3$, and integrating on $[0, u]$:

$$(*) \quad u = \int_{\wp(u)}^{+\infty} \frac{ds}{\sqrt{4s^3 - g_2s - g_3}}$$



- $\widehat{G}_{1\mathcal{E}_L}$ is a rotation with $\Theta(L) \in [0, \frac{1}{2}]$, and $\widehat{G}_{1\mathcal{E}_L}(\widehat{V}) = \widehat{V} + \widehat{H} = \widehat{H}$
- \widehat{H} has negative ordinate \Rightarrow is given by $u = 2\omega_1\Theta(L)$ and its abscissa is $X(L) = \wp(2\omega_1\Theta(L))$. Hence from (*):

$$2\omega_1\Theta(L) = \int_{X(L)}^{+\infty} \frac{ds}{\sqrt{4s^3 - g_2s - g_3}} \quad \Rightarrow \quad e_1 = \wp(\omega_1) \quad \Rightarrow \quad 2\Theta(L) = \frac{\int_{X(L)}^{+\infty} \frac{ds}{\sqrt{4s^3 - g_2s - g_3}}}{\int_{e_1}^{+\infty} \frac{ds}{\sqrt{4s^3 - g_2s - g_3}}} \sim \frac{4}{5}$$

References.

- Bastien, Rogalski; 2004. Global behavior of the solutions of Lyness' difference equation $u_{n+2}u_n = u_{n+1} + a$, *JDEA* 10.
- Cima, Gasull, Mañosa; 2007. Dynamics of the third order Lyness difference equation. *JDEA* 13.
- Feuer, Janowski, Ladas; 1996. Invariants for some rational recursive sequence with periodic coefficients, *JDEA* 2.
- Janowski, Kulenović, Nurkanović; 2007. Stability of the k th order Lyness' equation with period- k coefficient, *Int. J. Bifurcations & Chaos* 17.
- Jogia, Roberts, Vivaldi; 2006. An algebraic geometric approach to integrable maps of the plane, *J. Physics A* 39 (2006).
- Quispel, Roberts, Thompson; 1988-1989. Integrable mappings and soliton equations (II). *Phys. Lett. A* 126. and *Phys. D* 34.

Other Literature

- Bastien, Rogalski; 2007. On algebraic difference equations $u_{n+2} + u_n = \psi(u_{n+1})$ in \mathbb{R} related to a family of elliptic quartics in the plane, *J. Math. Anal. Appl.* 326.
- Beukers, Cushman; 1998. Zeeman's monotonicity conjecture, *J. Differential Equations* 143.
- Cima, Gasull, Mañosa; 2012a. On 2- and 3- periodic Lyness difference equations. *JDEA* 18.
- Cima, Gasull, Mañosa; 2012b. Integrability and non-integrability of periodic non-autonomous Lyness recurrences. arXiv:1012.4925v2 [math.DS]
- Cima, Zafar; 2012. Integrability and algebraic entropy of k -periodic non-autonomous Lyness recurrences. Preprint.
- Kulenović, Nurkanović; 2004. *Stability of Lyness' equation with period-three coefficient*, *Radovi Matematički* 12.
- Zeeman; 1996. Geometric unfolding of a difference equation. Unpublished paper.

THANK YOU!