

Hopf-Hopf and Hopf-Pitchfork bifurcations in coupled systems

Fátima Drubi, Santiago Ibáñez and Diego Noriega





February 6, 2023

- Coupled oscillatory systems: Isolated systems undergo a Hopf bifurcation
- Additional degeneracies may lead to Hopf-Hopf and Hopf-Pitchfork type bifurcations
- Likely, they become germs of complex bifurcation diagrams
- Particular case: Coupled neuron models

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(Figure from Quasar Jarosz at English Wikipedia (edited))



diffusively coupled





Isolated Oscillations

Dynamical Complexity?

Two identical systems...

diffusively coupled





Isolated Oscillations

Dynamical Complexity?

The Brusselator model:

$$\begin{cases} x' = A - (B+1)x + x^2 y \\ y' = Bx - x^2 y \end{cases}$$

with A and B positive.

The Coupled Brusselator System:

$$\begin{cases} x_1' = A - (B+1)x_1 + x_1^2 y_1 + \lambda_1 (x_2 - x_1) \\ y_1' = Bx_1 - x_1^2 y_1 + \lambda_2 (y_2 - y_1) \\ x_2' = A - (B+1)x_2 + x_2^2 y_2 + \lambda_1 (x_1 - x_2) \\ y_2' = Bx_2 - x_2^2 y_2 + \lambda_2 (y_1 - y_2) \end{cases}$$

where λ_1 and λ_2 are not negative.

Supercritical Hopf bifurcation



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Projection of a strange attractor:



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Interaction between two systems with *simple dynamics* can develop more *dynamical complexity* (e.g. chaos).

- [2] F. Drubi, S. Ibáñez, J. A. Rodríguez, J. Differential Equations 239 (2007).
- [3] F. Drubi, S. Ibáñez, J. A. Rodríguez, Bull. Belg. Math. Soc. Simon Stevin 15 (2008)
- [4] F. Drubi et al., Chaos, Solitons & Fractals (2023).

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where λ_1 and λ_2 are not negative.

In any generic unfolding X_{μ} , with $\mu \in \mathbb{R}^{n}$, of an *n*-dimensional nilpotent singularity of codimension *n*, there exist two bifurcation curves of (n - 1)-dimensional nilpotent singularities of codimension n - 1 which are generically unfolded by X_{μ} .

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3

The dynamics of coupled system in the synchronization plane

$$S = \{x_1 = x_2, y_1 = y_2\}$$

are those of the isolated system, i.e., there is a Hopf bifurcation in S.

The simplest dynamics expected by a transversal plane S^* are:

- a zero eigenvalue in its linear part that will lead to a **Hopf-Pitchfork bifurcation** of codimension two or higher.
- a pair of imaginary eigenvalues in its linear part that will lead to a Hopf-Hopf bifurcation of codimension two or higher.

Both bifurcations are *partially* studied in the literature.

- [6] J. Guckenheimer, P. Holmes, Nonlinear Oscillations, Dynamical Systems, and Bifurcations of Vector Fields, Springer (1983).
- [7] Y. Kuznetsov, Elements of applied bifurcation theory, 2ed, Springer (1998).

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The Coupled Brusselator System



[5] F. Drubi, S. Ibáñez, J. A. Rodríguez, *Physica D: Nonlinear Phenomena 240* (2011).

The Coupled Fitzhugh-Nagumo Oscillators

The FitzHugh-Nagumo (FN) model is raised from a traslation of Van der Pol's equation for a relaxation oscillator:

$$\begin{cases} x' = c \left(y + x - \frac{1}{3}x^3 + l \right) \\ y' = -\frac{1}{c} \left(x - a + b y \right) \end{cases}$$

with *x* the neuron **membrane potential** and *y* a **recovery variable**.

The **action potential** *I* corresponds to an external stimulus. Assuming

$$0 < b < 1$$
, $c > 0$, $1 - 2b/3 < a < 1$, and $b < c^2$

the system has a unique **atractor** (a **resting state**) when *I* = 0 and a **Hopf bifurcation** on the parameter *I*.

[8] R. Fitzhugh, *Biophysical Journal 1* (1961).

[9] M. Kawato, M. Sokabe, R. Suzuki, *Biological Cybernetics 34* (1979).

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Two FN neurons interacting **symmetrically** by a **linear coupling** is modelled by the 4-dimensional system

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with $\alpha_1 \leq 0$.

The plane $S = \{x_1 = x_2, y_1 = y_2\}$ is invariant by the flow.

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with lineal diffusion parameteres $\alpha_1, \alpha_2, \varepsilon_1, \varepsilon_2 \in \mathbb{R}$.

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The Coupled Fitzhugh-Nagumo Oscillators

A change of coordinates:

$$u_1 = \frac{1}{2}(x_1 - x_2), \quad v_1 = \frac{1}{2}(y_1 - y_2), \quad u_2 = \frac{1}{2}(x_1 + x_2), \quad v_2 = \frac{1}{2}(y_1 + y_2)$$

provides the equivalent equations

$$\begin{cases} u_1' = c \left(v_1 + u_1 - \frac{1}{3}u_1^3 - u_1u_2^2 \right) - (2\alpha_1 + \varepsilon_1)u_1 \\ v_1' = -\frac{1}{c}(u_1 + bv_1) - (2\alpha_2 + \varepsilon_2)v_1, \\ u_2' = c \left(v_2 + u_2 - \frac{1}{3}u_2^3 - u_1^2u_2 \right) + \varepsilon_1u_1, \\ v_2' = -\frac{1}{c}(u_2 - a + bv_2) + \varepsilon_2v_1. \end{cases}$$

with the invariant (**synchronization**) plane $\overline{S} = \{u_1 = v_1 = 0\}$. We will denote as \overline{S}^* to a plane transverse to \overline{S} .

The Coupled Fitzhugh-Nagumo Oscillators

Linear Analysis at the equilibrium on the invariant plane

The equilibrium on the invariant plane \overline{S} is of the form $(0, 0, p_2, q_2)$, where p_2 and q_2 fulfill the identities:

$$\frac{a}{b} + \left(1 - \frac{1}{b}\right)p_2 - \frac{1}{3}p_2^3 = 0$$
 and $q_2 = \frac{a - p_2}{b}$.

The linear part at $(0, 0, p_2, q_2)$ is given by

$$\begin{pmatrix} c\,(1-p_2^2)-(2\alpha_1+\varepsilon_1) & c & 0 & 0\\ -\frac{1}{c} & -\frac{b}{c}-(2\alpha_2+\varepsilon_2) & 0 & 0\\ \varepsilon_1 & 0 & c\,(1-p_2^2) & c\\ 0 & \varepsilon_2 & -\frac{1}{c} & -\frac{b}{c} \end{pmatrix}.$$

The Coupled Fitzhugh-Nagumo Oscillators

Elementary Local Bifurcations

$\mathcal{H}_{\overline{S}}$	p_2^2 = 1 $-b/c^2$, $b < c$
$\mathcal{P}_{\overline{\mathcal{S}}^*}$	$p_{2}^{2} = 1 - 1/(b + c(2\alpha_{2} + \varepsilon_{2})) - (2\alpha_{1} + \varepsilon_{1})/c$
$\mathcal{H}_{\overline{\mathcal{S}}^*}$	$p_2^2 = 1 - \frac{1}{c} \left(b/c + 2\alpha_1 + 2\alpha_2 + \varepsilon_1 + \varepsilon_2 \right)$ $1 - \left(b/c - (2\alpha_1 + \varepsilon_1) \right) \left(b/c + 2\alpha_2 + \varepsilon_2 \right) > 0$
$\mathcal{HP} = \mathcal{H}_{\overline{S}} \cap \mathcal{P}_{\overline{S}^*}$	$\begin{array}{l} p_2^2 = 1 - b/c^2, b < c \\ 1/(b + c(2\alpha_2 + \varepsilon_2)) + (2\alpha_1 + \varepsilon_1)/c = b/c^2 \\ \end{array}$ Hyperbolic eigenvalue: $-(2\alpha_1 + 2\alpha_2 + \varepsilon_1 + \varepsilon_2)$
$\mathcal{HH}=\mathcal{H}_{\overline{S}}\cap\mathcal{H}_{\overline{S}^*}$	$p_{2}^{2} = 1 - b/c^{2}, 2\alpha_{1} + 2\alpha_{2} + \varepsilon_{1} + \varepsilon_{2} = 0$ $(b/c + 2\alpha_{2} + \varepsilon_{2})^{2} < 1, b < c$

The Coupled Fitzhugh-Nagumo Oscillators

Resonance phenomena

In the Hopf-Hopf case, the eigenvalues are:

$$\pm i\sqrt{1-\left(rac{b}{c}-2lpha_1-arepsilon_1
ight)^2}$$
 and $\pm i\sqrt{1-rac{b^2}{c^2}}.$

If $2\alpha_1 + \epsilon_1 = 0$, the system will show a resonance 1:1. Hence, non-resonant bifurcations of codimension 2 can be unfolded.

As it occurs in **S. A. Campbell and M. Waite** (2001), where $\alpha_2 = \varepsilon_2 = 0$ and $\alpha_1 \le 0$.

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The Coupled Fitzhugh-Nagumo Oscillators

Bifurcation Diagram with AUTO: b = 0.4, c = 2, and $\alpha_2 = \varepsilon_2 = 0$



The Coupled Fitzhugh-Nagumo Oscillators





The Coupled Fitzhugh-Nagumo Oscillators





The Coupled Fitzhugh-Nagumo Oscillators





The Coupled Fitzhugh-Nagumo Oscillators

Bifurcation Diagram with AUTO: b = 0.4, c = 2, and $\alpha_2 = 0.01$



The Coupled Fitzhugh-Nagumo Oscillators



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The Coupled Fitzhugh-Nagumo Oscillators

Continuation of a periodic orbit ($\alpha_2 = 0.01$):



The Coupled Fitzhugh-Nagumo Oscillators

Modulus of eigenvalues (maximum) ($\alpha_2 = 0.01$):



The Coupled Fitzhugh-Nagumo Oscillators

 $a = 0.7, \alpha_1 = 0.072926, PO$



The Coupled Fitzhugh-Nagumo Oscillators

 $a = 0.7, \alpha_1 = 0.065822, 2$ -Torus



The Coupled Fitzhugh-Nagumo Oscillators

 $a = 0.7, \alpha_1 = 0.062702$, Chaotic Attractor?



Open Questions

- Is the case VIa associated with Hopf-Hopf singularities also feasible?
- What other cases can be obtained from resonance?
- Can all Hopf-Hopf types of codimension 2 be obtained?
- Can a classification of the Hopf-Pitchfork be provided?



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Thank you for your attention!