## P4 and desingularization of vector fields in the plane

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- P4 = Planar Polynomial Phase Portraits
- ▶ implemented by C Herssens, J C Artes, J Llibre, F Dumortier
- originally worked for unix with reduce
- Program ported to Qt (windows/unix/mac) with maple by PDM
- ▶ P5 = Piecewise P4

Workings of P4 is based on the book Qualitative Theory of Planar Differential Systems by Dumortier, LLibre and Artes.



$$\begin{cases} \dot{x} = P(x, y) \\ \dot{y} = Q(x, y) \end{cases}$$

Goal: qualitative study of dynamics, disregarding time-related features. This means looking at the phase portrait Theoretics:

- Poincare-Bendixson, so no chaos
- finite number of singular points when reduced
- study at infinity possible
- singular points have a finite number of sectors (parabolic, hyperbolic, elliptic)
- Separatrix skeleton can be drawn (problem of homoclinic and heteroclinic connections)
- Limit cycles may or may not be present

More than any phase portrait drawing program that one can easily find online!



$$\begin{cases} \dot{x} = P(x, y) \\ \dot{y} = Q(x, y) \end{cases}$$

Step 1: Eliminating GCF

This is done using Maple. In the sequel we will assume the GCF has been eliminated.

Step 2: Finding the isolated singular points. Some of them are evaluated algebraically some numerically, but all computations are done with real roots.

Step 3: behaviour at infinity Consider  $S^2 = \{X^2 + Y^2 + Z^2 = 1\}$ , and define  $\Delta(x, y) = \sqrt{1 + x^2 + y^2}$ ,

$$f^{\pm}(x,y) = \pm \left(rac{x}{\Delta},rac{y}{\Delta},rac{1}{\Delta}
ight) = (X,Y,Z)$$



$$\implies$$
 vf is defined on  $S^2$  outside equator

How to extend to the equator? Consider three charts

$$\phi_1(X, Y, Z) = \left(\frac{Y}{X}, \frac{Z}{X}\right) = (u, v)$$
  
$$\phi_2(X, Y, Z) = \left(\frac{X}{Y}, \frac{Z}{Y}\right)$$
  
$$\phi_3(X, Y, Z) = \left(\frac{X}{Z}, \frac{Y}{Z}\right) = (x, y)$$

Then define the vector field using the relation

$$(u, v) = (\phi_1 \circ \phi_3^{-1})(x, y)$$
  
=  $(y/x, 1/x)$ 

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The equator  $\{v = 0\}$  corresponds to infinity in the  $U_3$  chart.

$$(u, v) = (\phi_1 \circ \phi_3^{-1})(x, y) = (y/x, 1/x) \implies (x, y) = (1/v, u/v)$$
  
Chart  $U_2$ :

$$(u,v) = (\phi_2 \circ \phi_3^{-1})(x,y) = (x/y, 1/y) \implies (x,y) = (u/v, 1/v)$$

They can be joint by 1 formula:

$$(x,y) = \left(\frac{\cos\theta}{v}, \frac{\sin\theta}{v}\right)$$

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Chart  $U_1$ :

$$\begin{cases} \dot{x} = P(x, y) \\ \dot{y} = Q(x, y) \end{cases}$$

goes to

$$\begin{cases} \dot{u} = -uP(1/v, u/v) + Q(1/v, u/v) \\ \dot{v} = -vP(1/v, u/v) \end{cases}$$

and after multiplication to

$$\begin{cases} \dot{u} = v^{d} \left( -uP(1/v, u/v) + Q(1/v, u/v) \right) \\ \dot{v} = -v^{d+1}P(1/v, u/v) \end{cases}$$

where *d* is the degree of the polynomials *P*, *Q*. The result is again a polynomial vector field. At  $\{v = 0\}$ :

$$\begin{cases} \dot{u} = -uP_d(1, u) + Q_d(1, u) \\ \dot{v} = 0 \end{cases}$$

Equator is invariant with a well-defined dynamics on it!

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## P4 shows a view of the sphere from the top:



Poincaré compactification:

$$(x,y) = \left(\frac{\cos\theta}{v}, \frac{\sin\theta}{v}\right)$$

Poincaré-Lyapunov compactification:

$$(x,y) = \left(\frac{\cos\theta}{v^{\alpha}}, \frac{\sin\theta}{v^{\beta}}\right)$$

Same idea buth with weights  $(\alpha, \beta)$  (and with a bit more complicated inverted formula)

Step 4: Local study of singular points

$$\begin{cases} \dot{x} = P(x, y) \\ \dot{y} = Q(x, y) \end{cases}$$

Suppose

$$P(x_0, y_0) = Q(x_0, y_0) = 0.$$

Define the jacobian

$$M = \begin{pmatrix} \frac{\partial P}{\partial x}(x_0, y_0) & \frac{\partial P}{\partial y}(x_0, y_0) \\ \frac{\partial Q}{\partial x}(x_0, y_0) & \frac{\partial Q}{\partial y}(x_0, y_0) \end{pmatrix}$$

and consider the linearized equation

$$\left(\begin{array}{c} \dot{x} \\ \dot{y} \end{array}\right) = M \left(\begin{array}{c} x - x_0 \\ y - y_0 \end{array}\right)$$

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Several cases:

- 1. Saddle (eigenvalues  $\lambda, \mu$  opposite sign)
- 2. Node (eigenvalues  $\lambda, \mu$  same sign and nonzero)
- 3. Focus (eigenvalues  $\alpha \pm i\beta$ ,  $\alpha \neq 0$ ,  $\beta \neq 0$ )
- 4. Center (eigenvalues  $\pm i\beta$ ,  $\beta \neq 0$ )
- 5. Semi-elementary (eigenvalues  $\lambda$ , 0 with  $\lambda \neq 0$ )
- 6. nilpotent or degenerate (eigenvalues 0,0)

For case 1: we compute invariant manifolds tangent to eigenspace of  $\lambda$  resp.  $\mu.$ 

For cases 4,5,6 we need information from the nonlinear part to determine the type further

Case 4: Lyapunov constants (see talk of Joan Torregrosa). P4 uses a method of Gasull & Torregrosa

Case 5: there exists a smooth 1-dim center manifold which is a graph y = h(x) or x = k(y). Reduction of the dynamics to the center manifold leads to determination of type.

Case 6: desingularization

Consider a singular point at the origin (0,0). We use

$$(x, y) = (r \cos \theta, r \sin \theta) = (r\overline{x}, r\overline{y}).$$

and use  $(r, \theta)$  as new coordinates. Near  $\theta = 0$  we use  $\sin \theta \approx \theta$  and  $\cos \theta \approx 0$ , so

$$(x,y)=(r,r\theta)$$

Better:

$$(x,y) = (r,r\overline{y})$$
 "chart  $\overline{x} = 1$ "

Near  $heta=\pi/2$  we have  $\sin heta pprox 1$  and  $\cos heta pprox heta -\pi/2$ , so

$$(x,y) = (r(\theta - \pi/2), r)$$

Better:

$$(x,y) = (r\overline{x},r)$$
 "chart  $\overline{y} = 1$ "

Instead of using  $(r, \theta)$  we use the charts.

Example:

$$\begin{cases} \dot{x} = x^2 - 2xy \\ \dot{y} = y^2 - xy \end{cases}$$

Leads to

$$\begin{cases} \dot{r} = r(\cos^3\theta - 2\cos^2\theta\sin\theta + \dots) + O(r^2) \\ \dot{\theta} = \cos\theta\sin\theta(3\sin\theta - 2\cos\theta) + O(r) \end{cases}$$

Seems somewhat complicated trigonometry but is in fact not so hard



It is better to use the charts instead of  $(r, \theta)$ :

$$egin{aligned} & (x,y) = (r,r\overline{y}) & \ & ext{``chart } \overline{x} = 1'' \ & & & \\ & \dot{x} &= x^2 - 2xy \ & \dot{y} &= y^2 - xy \end{aligned}$$

Leads to

$$\begin{cases} \dot{r} = r(1-2\overline{y}) \\ \dot{\overline{y}} = 3\overline{y}^2 - 2\overline{y} \end{cases}$$

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 $\implies$  polynomial character is retained. Of course to get information on the full circle we need to complement with additional charts. Sometimes more than one blow-up is necessary:



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Theorem: any singular point of an analytic planar vector field can be blown up after a finite number of blowups so that on the blow-up locus only elementary or semi-elementary singular points are found



For each of these (semi)elementary points one can compute separatrices.

 $\implies$  for any singular point there is an algorithm to divide the neighbourhood in sectors (hyperbolic, eliptic, parabolic) and to compute the type of the singular point.

P4 actually implements Quasi-homogeneous blow-up

$$(x,y) = (r^{\alpha} \cos \theta, r^{\beta} \sin \theta) = (r^{\alpha} \overline{x}, r^{\beta} \overline{y}).$$

How to choose the weights  $(\alpha, \beta)$ ? Let

$$\dot{x} = P(x, y) = \sum a_{ij} x^{i} y^{j}, \dot{y} = Q(x, y) = \sum b_{ij} x^{i} y^{j}$$
$$S = \{(i - 1, j) : a_{ij} \neq 0\} \cup \{(i, j - 1) : b_{ij} \neq 0\}$$

The newton polygon is the convex hull of the set

$$\mathcal{P} = \cup_{(r,s)\in S} \{ (r',s') \colon r' \geq r, s' \geq s \}.$$

One of the borders of the Newton polygon is a straight line with equation

$$r\alpha + s\beta = m$$

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then  $(\alpha, \beta)$  is a suitable choice

Lemma: if we proceed this way, then after blowing up, the north and south poles are either nonsingular or (semi)elementary

 $\implies$  iterated blow-ups are only necessary in the horizontal directions.

This reduces the computational work.

Conclusion: besides determining homoclinic, heteroclinic connections and limit cycles, P4 offers a full global study of planar vector fields.

P5: same thing but with piecewise polynomial systems, defined in regions by algebraic inequalities

Possible extensions to P4/P5:

- computing saddle quantities
- alternative algorithms for numerical integration
- beter sewing in P5
- period computation, computing abelian integrals, Melnikov integrals, ...

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